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## **Novel 3-D PtS-like Tetrazolate-Bridged Manganese(II) Complex Exhibiting Spin-Canted Antiferromagnetism and Field-Induced Spin-Flop Transition**

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Two novel manganese(II) tetrazolate coordination polymers,  $Mn_3(Hbta)_4(\mu_ZOH_2O) \cdot 2H_2O$  (1) and  $Mn(bta)(H_2O)$ (**2**), were obtained by the hydrothermal reaction of MnCl2 · 4H2O with *N*,*N*-bis[1(2)*H*-tetrazol-5-yl]amine (H2bta) in different pH values. Compound 1 exhibits a 2-D (4,4) layer structure, which is built from binuclear  $Mn_2(\mu_2$ -OH)<sub>2</sub> subunits and  $\mu_3$ -bridging Hbta<sup>-</sup> linkers. Compound 2 is a helical framework, and its 3-D PtS-like net is constructed by  $Mn_2(bta)_2$  dimers and  $\mu_5$ -bridging bta<sup>2-</sup> linkers. Magnetic measurements reveal that compound 1 displays an antiferromagnetic ordering, while compound **2** exhibits the coexistence of a spin-canted antiferromagnetic coupling and a field-induced spin-flop transition.

### **Introduction**

Considerable current interest is focused on the preparation and study of molecular-based magnetic materials because of their great value in investigating complex magnetostructural relationships and some fundamentally magnetic behaviors.<sup>1</sup> In particular, molecular magnetism with spontaneous magnetization, such as spin canting and metamagnetism, is continously pursued for its potential application as molecular magnets.<sup>2</sup> However, the design and preparation of such molecular magnetic materials are still a challenging topic. The common synthetic approach of molecular magnetic materials is to incorporate paramagnetic centers with suitable bridging ligands to fabricate polynuclear complexes or extended networks that can give rise to magnetic exchange

coupling. Besides the nature of the metal centers, the nature and the linking modes of the bridging ligand also play a crucial role in the whole magnetic system; therefore, the search and exploration of an appropriate ligand is a longstanding task. $3$ 

As bridging ligands, five-membered heterocyclic azoles (pyrazoles, imidazoles, triazoles, tetrazoles, etc.) and their derivatives have drawn much attention because of their strong *σ*-electron donors to metals and their diverse coordination modes, forming numerous multidimensional coordination complexes with intriguing topologies.4 They have also shown to be potential candidates to construct molecular magnets, exhibiting various bulk magnetic behaviors such as antiferromagnetism, ferromagnetism, spin crossover, spin canting,

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ferrimagnetism, and so on.<sup>5</sup> With respect to tetrazolate ligands, quite a few magnetical systems including iron $(II)$ ,<sup>6</sup> copper(II),<sup>7</sup> cobalt(II),<sup>8</sup> and iron(III) and lanthanide(III)<sup>9</sup> have been characterized. However, to the best of our knowledge, up to now the magnetic investigation on tetrazole-bridged manganese(II) complexes is very rare<sup>10</sup> because only a few polynuclear or multidimensional manganese(II) complexes of tetrazole-containing ligands have been observed, and in some cases, the magnetic studies have not been described.<sup>11</sup> Although numerous tetrazolate-containing complexes with spin crossover, antiferromagnetism, ferromagnetism, and

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ferrimagnetism have been found,  $6-10$  there are only two examples displaying spin-canted behavior, $12$  and in particular those with spin-flop behavior have not yet been reported.

The *N*,*N*-bis[ $1(2)$ *H*-tetrazol-5-yl]amine (H<sub>2</sub>bta) ligand as an important tetrazaolate derivative remains mostly unexplored except for the recently reported several mononuclear, <sup>13</sup> 1-D,<sup>7i</sup> and 2-D<sup>10</sup> coordination polymers. We choose the H<sub>2</sub>bta ligand in the present work on the basis of the following considerations: (1) it contains nine-coordinate atoms originating from double tetrazole rings and an amine group, promising rich and versatile coordination modes to achieve diverse structures with interesting topology and unexpected magnetic behaviors; (2) it possesses three reversible diprotonated, monoprotonated, and deprotonated types of H2bta,  $Hbta^{-}$ , and  $bta^{2-}$ , implying diverse acidity-dependent coordination modes, and (3) it is a flexible bridging ligand, favoring the formation of a helical structure. In this paper, two unprecedented manganese(II) tetrazolate complexes,  $Mn_3(Hbta)_4(\mu_2-OH)_2(H_2O)\cdot 2H_2O (1)$  and  $Mn(bta)(H_2O) (2)$ , were synthesized under hydrothermal conditions. Notably, the different acidity-dependent types and coordination modes of the ligand lead to distinct dimensions and architectures of the two complexes. Magnetic studies show that compound **1** exhibits an antiferromagnetical ordering with  $T_N = 4.5$  K and compound **2** shows the simultaneous presence of a spincanted antiferromagnetism and a field-induced spin-flop transition.

#### **Experimental Section**

**Materials and Instruments.** All chemicals were used as purchased without further purification. H2bta was synthesized according to the literature.<sup>71</sup> Powder X-ray diffraction (PXRD) patterns were obtained from a Rigaku DMAX2500 powder diffractometer at 40 kV and 100 mA using Cu K $\alpha$  ( $\lambda$  = 1.540 56 Å), with a scan speed of 0.375 s/step and a step size of 0.05°. The simulated PXRD patterns were calculated using single-crystal X-ray diffraction data and processed by the free *Mercury* v1.4 program provided by the Cambridge Crystallographic Data Centre. Magnetic susceptibilities were acquired by using a Quantum Design PPMS model 6000 magnetometer. Elemental analyses of C, H, and N were carried out with an Elementar Vario MICRO. The Fourier transform infrared (FT-IR) spectra were measured on a Nicolet Magna 750 FT-IR spectrophotometer in the region of  $4000-400$  cm<sup>-1</sup> using KBr pellets as bases. Thermogravimetric analyses (TGA) were investigated on a Netzsch Sta449C thermoanalyzer under a  $N_2$ atmosphere with a ramp rate of 10 °C/min.

**Preparations of 1 and 2.** A mixture of  $MnCl_2 \cdot 4H_2O$  (85 mg, 0.8 mmol) and  $H_2$ bta (84 mg, 0.5 mmol) in 10 mL of deionized water with a starting value of  $pH = 3$  was sealed into a 25-mL poly(tetrafluoroethylene)-lined stainless steel container under autogenous pressure, kept at 160 °C for 3 days, and then cooled to room temperature at a rate of 2.7 °C/h. Pale-red crystals of **1** were filtered off from the solution (yield 40%, based on Mn), then washed with water, and dried in air. Anal. Calcd for  $C_8H_{18}Mn_3N_{36}O_6$ : C, 10.93; H, 2.06; N, 57.34. Found: C, 11.08; H, 1.91; N, 56.61. IR (KBr, cm-1): *ν* 3292 (s), 3041 (s), 2916 (s), 2789 (s), 2629 (m),

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**Table 1.** Crystal Data for **1** and **2**



1871 (w), 1659 (vs), 1630 (vs), 1508 (vs), 1431 (m), 1319 (m), 1234 (w), 1146 (w), 1041 (m), 991 (w), 895 (w), 854 (w), 798 (m), 731 (m), 428 (w).

Colorless crystals of **2** in 20% yield (based on Mn) were obtained by a method similar to that of **1** except that 9 mL of deionized water and 1 mL of  $25\% \text{ NH}_3 \cdot \text{H}_2\text{O}$  were used as the solvents with a starting value of  $pH = 11$ , and then the mixture was filtered off from the solution, washed with water, and dried in air. Anal. Calcd for C2H3MnN9O: C, 10.72; H, 1.35; N, 56.26. Found: C, 10.73; H, 1.33; N, 55.46. IR (KBr, cm-1): *ν* 3278 (s), 3167 (m), 3089 (s), 2934 (s), 2185 (w), 1633 (s), 1551(s), 1514 (s), 1411 (m), 1336 (w), 1271 (w), 1146 (m), 1034 (w), 864 (m), 810 (s), 721 (m), 621 (m), 434 (w).

**Single-Crystal Structure Determination.** Single crystals were mounted on a Rigaku Mercury CCD diffractometer equipped with graphite-monochromated Mo K $\alpha$  radiation ( $\lambda = 0.710$  73 Å) using an *ω*-scan technique at 293 K. The intensity data sets were reduced by the *CrystalClear* software.<sup>14</sup> The structures were solved by direct methods with the Siemens *SHELXTL* version 5 package of crystallographic software<sup>15</sup> and refined by full-matrix least-squares techniques. Non-H atoms were located by difference Fourier maps and subjected to anisotropic refinement. H atoms were added according to theoretical models. Crystallographic data and structural refinements and selected bond distances and angles for **1** and **2** are listed in Tables 1 and 2, respectively.

**TGA for 1 and 2.** TGA of **1** show 21.2% weight loss in the range of 30-<sup>200</sup> °C (Figure S1a in the Supporting Information), in agreement with the removal of solvent and coordinated water molecules per cell (12.65% calcd). Further weight losses from 200 to 500 °C correspond to decomposition of the network. The TGA curve of **2** reveals that the framework was stable between 30 and 360 °C and then weight losses occurred from 360 to 864 °C due to decomposition of the network (Figure S1b in the Supporting Information).

**Table 2.** Selected Bond Lengths and Angles for **1** and **2***<sup>a</sup>*

complex <sub>1</sub>		complex <sub>2</sub>	
Bond Length			
$Mn(1) - O(1W)$	2.136(2)	$Mn(1) - O(1W)$	2.192(2)
$Mn(2)-O(1)$	2.359(2)	$Mn(1)-N(1A)$	2.276(2)
$Mn(2)-N(9)$	2.172(2)	$Mn(1) - N(2B)$	2.302(2)
$Mn(1)-N(3)$	2.302(2)	$Mn(1)-N(4)$	2.297(2)
$Mn(2)-N(1)$	2.261(2)	$Mn(1)-N(7C)$	2.316(2)
		$Mn(1)-N(9A)$	2.245(2)
Angle			
$O(1W) - Mn(1) - O(1WA)$	180.00(1)	$O(1W) - Mn(1) - N(9A)$ 110.56(8)	
$N(9) - Mn(2) - N(1)$	78.87(6)	$N(9A) \# 1 - Mn(1) - N(4) 84.04(8)$	
$O(1W) - Mn(1) - N(3)$	90.84(7)	$N(6)-N(7)-Mn(1C)$	116.39(16)
$N(1) - Mn(2) - O(1)$	97.22(8)	$N(1A) - Mn(1) - N(2B)$	99.73(7)
$N(9) - Mn(2) - N(1D)$	103.83(6)	$N(2B) - Mn(1) - N(7C)$	76.21(8)
$N(3)-Mn(1)-N(3A)$	180.00(5)	$N(3)-N(4)-Mn(1)$	116.56(14)
$N(3B) - Mn(1) - N(3A)$	94.16(9)	$N(2)-N(1)-Mn(1D)$	128.24(15)
$N(1) - Mn(2) - O(1D)$	159.04(4)	$N(4) - Mn(1) - N(7C)$	100.10(8)
$O(1D) - Mn(2) - O(1)$	73.13(10)	$O(1W) - Mn(1) - N(7C)$	158.06(8)
$N(9) - Mn(2) - O(1)$	93.63(4)	$O(1W) - Mn(1) - N(1A)$ 85.11(8)	
$Mn(2E) - O(1) - Mn(2)$		$106.87(10) N(9A) - Mn(1) - N(1A)$	78.95(7)
$C(2)-N(5)-C(1)$		$124.85(14) N(9A) - Mn(1) - N(7C)$	91.38(8)
		$N(1)-N(2)-Mn(1E)$	131.98(16)
		$C(1)-N(5)-C(2)$	123.5(2)

*a* Symmetry codes for **1**: (A)  $-x$ ,  $-y$  + 1,  $-z$ ; (B) *x*,*y*,  $-z$ ; (C)  $-x$ ,  $-y$ <br>1 *z*: (D)  $-x$  + 1,  $-y$  + 1,  $z$ ; (E)  $-x$  + 1,  $y$  + 0  $-z$  + 1/<sub>2</sub>. Symmetry + 1, *z*; (D)  $-x + 1$ ,  $-y + 1$ , *z*; (E)  $-x + 1$ ,  $y + 0$ ,  $-z + \frac{1}{2}$ . Symmetry codes for 2: (A)  $x - \frac{1}{2}$ ,  $-y + \frac{3}{2}$ ,  $z + \frac{1}{2}$ ; (B)  $-x + \frac{1}{2}$ ,  $y - \frac{1}{2}$ ,  $-z + \frac{3}{2}$ ; codes for **2**: (A)  $x - {}^{1}/2$ ,  $-y + {}^{3}/2$ ,  $z + {}^{1}/2$ ; (B)  $-x + {}^{1}/2$ ,  $y - {}^{1}/2$ ,  $-z + {}^{3}/2$ ;<br>(C)  $-x - y + 1$ ,  $-z + 1$ ; (D)  $x + {}^{1}/2$ ,  $-y + {}^{3}/2$ ,  $z - {}^{1}/2$ ; (E)  $-x + {}^{1}/2$ , y (C)  $-x$ ,  $-y$  + 1,  $-z$  + 1; (D)  $x$  +  $\frac{1}{2}$ ,  $-y$  +  $\frac{3}{2}$ ,  $z$  -  $\frac{1}{2}$ ; (E)  $-x$  +  $\frac{1}{2}$ ,  $y$  +  $\frac{1}{2}$ ,  $-z$  +  $\frac{3}{2}$  $+$   $\frac{1}{2}$ ,  $-z$   $+$   $\frac{3}{2}$ .

#### **Results and Discussion**

**Syntheses.** As we know, in a special hydrothermal process, many factors can affect the formation and crystal growth of the product phases, such as the initial reactants, starting concentrations, solvents, pH values, reaction times, temperatures, etc. In our case, the pH value of the reaction solution is a key factor in the formation of different structural networks in **1** and **2** mainly resulting from the aciditydependent modes of the ligand. Under acidic conditions, the ligand is monodeprotonated and adopts the Hbta<sup>-</sup> type for **1**; on the contrary, it employs the  $bta^{2-}$  dideprotonated type for **2** under the basic environment. The PXRD patterns of the bulk products are in good agreement with the calculated patterns based on the single-crystal solution, indicating the phase purity of **1** and **2** (Figure S2 in the Supporting Information).

**Crystal Structures of 1 and 2.** A single-crystal X-ray analysis shows that the framework of **1** features a 2-D gridlike network. In the asymmetrical unit of **1** (Figure 1), there are two crystallographically independent Mn<sup>II</sup> ions, one Hbta<sup>-</sup> ligand, one hydroxyl group, one coordinated water molecule, and one solvated water molecule. The Mn1 atom is trans-coordinated by four N3 atoms from four Hbta<sup>-</sup> ligands on the equatorial plane, with  $Mn1-N3 = 2.302(2)$ Å and two terminal water molecules occupying the two axial sites with Mn1 $-O1W = 2.136(2)$  Å, yielding a compressed  $(4 + 2)$  octahedron. The Mn2 atom is cis-coordinated by two chelating Hbta<sup>-</sup> ligands with  $Mn2-N1 = 2.261(2)$  and  $Mn2-N9 = 2.172(3)$  Å and two hydroxyl groups with  $Mn2-O1 = 2.359(2)$  Å to complete a slightly distorted octahedron. In **1**, each of Hbta<sup>-</sup> ligand acts as a  $\mu_3$  linker, which coordinates to the Mn2 ion in a chelating fashion through N1 and N9 atoms from different tetrazole rings and

<sup>(14)</sup> *CrystalClear: Software User's Guide for the Rigaku R-Axis, and Mercury and Jupiter CCD Automated X-ray Imaging System*, version 1.35; Rigaku Molecular Structure Corp.: The Woodlands, TX, 2002.

<sup>(15)</sup> *SHELXTL Reference Manual*, version 5; Siemens Energy & Automation Inc.: Madison, WI, 1994.



Figure 1. ORTEP drawing with 50% thermal ellipsoids of the manganese coordination environment in **1**. Symmetry codes: (A)  $-x$ ,  $1 - y$ , *z*; (B) 1  $-x$ ,  $1-y$ ,  $z$ ; (C)  $x$ ,  $1-y$ ,  $\frac{1}{2}-z$ .

also bonds to a Mn1 atom via the N3 atom, presenting a  $similar -Mn-N1-N2-N3-Mn-$ end-to-end azide-bridged coordination mode to link Mn1 and Mn2  $[Mn1 \cdots Mn2]$ 6.487(4) Å]. The two tetrazole rings of each  $Hbta$ <sup>-</sup> ligand are almost coplanar with a small dihedral angle of  $4.23(6)^\circ$ ; moreover, the two ligated Hbta<sup>-</sup> ligands binding to the Mn2 atom are nearly perpendicular with a dihedral angle of 86.36(3)°. Two  $\mu_2$ -OH groups doubly bridge two equivalent Mn2 atoms to form a rhombus binuclear  $Mn_2(\mu_2\text{-}OH)_{2}$  unit with a  $Mn2 \cdots Mn2$  separation of 3.789(4) Å and a Mn2-O1-Mn2C bond angle of  $106.9(1)$ °. Each Mn<sub>2</sub>( $\mu$ <sub>2</sub>-OH)2(Hbta)4 dimer links four Mn1 centers, while each Mn1 center connects four  $Mn_2(\mu_2-OH)_2(Hbta)_4$  dimers, forming a novel 2-D (4,4) topological network in the *ac* plane (Figure 2). Notably, only five complexes<sup>16</sup> containing  $Mn_2(\mu_2-OH)_2$ units are observed in the Cambridge Structural Database, and all of them are discrete in structure. As a result, complex **1** is the first extended architecture constructed by  $Mn_2(\mu_2\text{-}OH)_2$ units. Moreover, taking into account the strong hydrogenbonding interactions [O1-H··· N7, O1W-H··· O2W, N5-H ··· N5, and N6-H ··· N6, with distances ranging from  $2.742(4)$  to  $2.867(6)$  Å and angles falling in the range of  $119.2(5)-180(2)$ °], complex 1 can be extended to a 3-D supramolecular architecture but without open channels because the adjacent layers are arranged in an  $-ABAB$ manner (Figure S3 in the Supporting Information).

Compound **2** crystallizes in the monoclinic space group *P*21/*c*, and its 3-D framework is constructed by binuclear  $Mn<sub>2</sub>(bta)<sub>2</sub>$  units. There is a unique Mn atom in the asymmetric unit. Each Mn ion adopts a distorted octahedral MnN<sub>5</sub>O geometry, with the equatorial plane consisting of four N atoms from three  $bta^{2-}$  ligands and the axial positions being occupied by one N atom from another  $bta^{2-}$  ligand and one O atom from the terminal water molecule (Figure 3). The bond lengths of Mn-N range from 2.245(2) to 2.316(2) Å and that of Mn-O is 2.192(2) Å, comparable to



**Figure 2.** 2-D layer of **1** viewed along the *b* axis (H atoms are omitted for clarity).



**Figure 3.** ORTEP view with 50% thermal ellipsoids of the  $Mn_2(bta)_2$  core in 2. Symmetry codes: (A)  $\frac{1}{2} - x$ ,  $\frac{-1}{2} - y$ ,  $\frac{3}{2} - z$ ; (B)  $-x$ ,  $1 - y$ ,  $1 - z$ <br> *z*; (C)  $\frac{-1}{2} - x$ ,  $\frac{3}{2} + y$ ,  $\frac{1}{2} + z$ ; (D)  $\frac{1}{2} - x$ ,  $\frac{1}{2} + y$ ,  $\frac{3}{2} - z$ ; (E)  $1 - x$ , 2 *z*; (C)  $-\frac{1}{2} - x$ ,  $\frac{3}{2} + y$ ,  $\frac{1}{2} + z$ ; (D)  $\frac{1}{2} - x$ ,  $\frac{1}{2} + y$ ,  $\frac{3}{2} - z$ ; (E)  $1 - x$ , 2  $- y$ ,  $1 - z$ .

those in the literature.<sup>17</sup> Interestingly, different from that in **1**, the H<sub>2</sub>bta ligand employs the dideprotoned bta<sup>2-</sup> type under basic conditions and behaves as a pentadentate coordination mode. In the mononuclear,  $^{13}$  1-D, $^{7i}$  and 2-D<sup>10</sup> structures, the  $bta^{2-}$  ligand acts as a bidentate, tridentate, or tetradentate linker, respectively; however, the pentadentate coordination mode of the  $bta^{2-}$  ligand exhibited in 2 is first observed.

In  $2$ , each bta<sup>2-</sup> ligand is bonded to four Mn atoms, with N1 and N9 atoms chelating the Mn1C center, the N2 atom binding the neighboring Mn1D atom, and N4 and N7 atoms connecting the other two Mn (Mn1 and Mn1B) centers, thus setting up two end-to-end azide-like linkages  $(-Mn-N^2-N-NA-Mn-$  and  $-Mn-N^7-N^8-N^9-Mn-$ ) in each

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**Figure 4.** Plots of the (a) left-handed and (b) right-handed chiral layers constructed by the two types of helical chains in **2**.

tetrazole ring. The dihedral angle between the two tetrazole rings is 19.64(6)°, which is larger than that in **1**. As shown in Figure 3, two Mn1 atoms are connected by two  $bta^{2-}$ ligands to construct a dimeric  $Mn_2(bta)_2$  unit, in which two  $MnC_2N_3$  six-membered rings are built through each bta<sup>2-</sup> ligand chelating one Mn1 atom and one  $Mn_2N_4$  sixmembered ring is constructed through two N1-containing tetrazole rings doubly bridging two Mn1 atoms. The similar  $M_2N_4$  six-membered rings were previously found in two tetrazolate-bridged zinc(II) complexes.<sup>18</sup> The Mn $\cdots$ Mn separations connected by double-bridging N1-containing tetrazole rings and by the above-mentioned end-to-end azidelike bridging mode are 4.295(3) and 6.243(1) Å, respectively.

Notably, the adjacent  $Mn_2(bta)_2$  units are interrelated by a 21-screw operation. To shed light on the helix and the structure of **<sup>2</sup>**, we first omit the N5-C1/C2 bonds in the  $bta<sup>2</sup>$  ligands and the Mn1-N1 bonds. As shown in Figure 4, each helical chain can be described as  $(-Mn-N-N (N-)$ <sub>n</sub>, and there are two types of helices ( $-Mn-N2-N3-$ N4-Mn- and -Mn-N7-N8-N9-Mn-) from the two tetrazole rings in each bta<sup>2-</sup> ligand. The pitch of the  $2<sub>1</sub>$  helix is about 9.711 Å along the *b* axis. The two types of helical chains with the same handness are interwoven via Mn atoms as hinges to extend into a 2-D chiral layer with left- or righthandness (parts a and b of Figure 4, respectively). As shown in Figure 5, these chiral layers are stacked alternately through Mn1-N1, N5-C1, and N5-C2 bonds to form a 3-D network of **2** with nonchirality. It is noteworthy that compound **2** represents the first bta-bridged complex featuring a 3-D structure. Taking into account each  $bta^{2-}$  ligand and the Mn center as four nodes, the framework of **2** exhibits a distorted 3-D PtS net with the Schläfli topological notation for 42 84 , in which the long vertex symbols for the fourconnecting-node  $bta^{2-}$  ligands and Mn centers are 4.4.8<sub>2</sub>.8<sub>2</sub>.8<sub>8</sub>.8<sub>8</sub> and 4.4.8<sub>7</sub>.8<sub>7</sub>.8<sub>7</sub>,8<sub>7</sub>, respectively.<sup>19</sup>

**Magnetic Properties.** The plots of the  $\chi_M$  and  $\chi_M T$  vs *T* and the fitting line under a 5 kOe external field for **1** are shown in Figure 6. The  $\chi_M T$  value at 300 K is 12.76 emu · K/ mol, being close to the spin-only value of  $13.12$  emu $\cdot$ K/



**Figure 5.** View the 3-D network of **2** along the *a* axis (pale red for righthanded layers and green for left-handed layers, and O atoms are omitted for clarity).



**Figure 6.** Temperature dependence of  $\chi_M$  and  $\chi_M T$  for 1 (inset: the d( $\chi_M T$ )/ d*T* vs *T* derivative curve).

mol expected for three magnetically isolated high-spin Mn<sup>II</sup> ions considering  $g = 2.00$ . During cooling, the value of  $\chi<sub>M</sub>T$ continues to slowly decrease until near 100 K; below this temperature, the  $\chi_M T$  value decreases sharply to 1.04 emu $\cdot$ K/ mol at 2 K, suggesting antiferromagnetic behavior. The appearance of a round peak in the  $\chi_M$  vs *T* curve indicates

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antiferromagnetic ordering. The magnetic susceptibility data in the temperature range of  $20-300$  K can be well fitted to the Curie-Weiss law expression with  $C = 13.15$  emu · K/ mol and  $\theta = -10$  K, which further confirms an overall antiferromagnetic interaction between the  $Mn^{\text{II}}$  ions. The Neel temperature,  $T_N = 4.5$  K, was determined from the sharp peak in  $d(\chi_M T)/dT$  (Figure 6, inset).

To quantitatively analyze the magnetic behavior of **1**, we assume that the magnetic coupling of **1** is dominated by the  $Mn_2(\mu_2\text{-}OH)_{2}$  core constituted by two Mn2 ions and the Mn1 ion is treated as a paramagnetic ion on the basis of the structure analysis, in which the separation of  $Mn2 \cdots Mn2$  $[3.789(4)$  Å connected by double OH<sup>-</sup> bridges is much shorter than that of Mn1 $\cdots$ Mn2 [6.487(4) Å] by three N atoms of the tetrazole ring. Accordingly, the molar magnetic susceptibility ( $\gamma$ ) of 1 can be described as  $\gamma = \gamma_C + (5/2)(7/2)$  $2)(Ng^2\beta/3KT)$ , in which  $\chi_C = (2Ng^2\beta^2/KT)(A/B)$  is for the M<sub>D</sub>, dimer<sup>20</sup> based on the isostropic spin Hamiltonian with  $Mn<sub>2</sub>$  dimer<sup>20</sup> based on the isostropic spin Hamiltonian with  $H = -2JS_1S_2$  and the latter is the contribution of a paramagnetic Mn1 ion. Moreover, considering the intermolecular interactions (*zJ*′), the magnetic data between 20 and 300 K can be well fitted using the following expression:

$$
\chi_{\rm M} = \chi/(1 - zJ' \chi N g^2 \beta^2)
$$
 (1)

$$
A = e^{2y} + 5e^{6y} + 14e^{12y} + 30e^{20y} + 55e^{30y}
$$
 (2)

$$
B = 1 + 3e^{2y} + 5e^{6y} + 7e^{12y} + 9e^{20y} + 11e^{30y}
$$
 (3)

where  $y = J/KT$ . The best-fitting parameters obtained are  $g = 1.98$ ,  $I = -1.36$ , cm<sup>-1</sup>, and  $zI' = -0.004$ , cm<sup>-1</sup>. The  $= 1.98, J = -1.36$  cm<sup>-1</sup>, and  $zJ' = -0.004$  cm<sup>-1</sup>. The agreement factor *R*, defined as  $\sum [(\chi_M T)_{obsd} - (\chi_M T)_{call}^2]^2$ <br> $\sum (\chi_M T)_{\text{c}}^2$  is equal to 6.0 × 10<sup>-4</sup>. The negative *I* value  $\sum(\chi_M T)_{obsd}^2$ , is equal to  $6.0 \times 10^{-4}$ . The negative *J* value reveals that dimeric Mn<sup>II</sup> spins connected by double  $\mu_2$ hydroxo bridging are antiferromagnetic coupling in **1**, which is consistent with most of doubly oxygen-bridged dimanganese(II) complexes.<sup>21,16b</sup>

At low-temperature region of  $2-20$  K in the low external field of 100 Oe, the curves of field-cooled (FC) and zerofield-cooled (ZFC) magnetizations (Figure 7) are identical, confirming an antiferromagnetic ordering. Furthermore, the real part  $(\chi')$  of the alternating current (ac) magnetic susceptibility in **1** has a maximum, and the imaginary part  $(\chi'')$  of the ac magnetic susceptibility is negligibly small (Figure S4a in the Supporting Information). No frequency dependence was observed in the ac measurements. This also supports the conclusion that compound **1** possesses a magnetic transition phase. At 2 K in the applied field of up to 70 kOe, the field dependence of magnetization increases



**Figure 7.** FC and ZFC ( $H = 100$  Oe) magnetization for **1**.



**Figure 8.** Temperature dependence of  $\chi_M$  and  $\chi_M T$  for 2.

almost linearly, reaching  $8.48 \text{ N} \beta$  at 70 kOe (Figure S5a in the Supporting Information), which is far from the saturation value of 15  $N\beta$  for three Mn<sup>II</sup> ions ( $S = \frac{5}{2}$ ). This behavior<br>is characteristic of ontiferromagnetic exchange counting is characteristic of antiferromagnetic exchange coupling.

The magnetic susceptibility of **2** is presented in Figure 8. The temperature dependence of the reciprocal susceptibility  $(\chi_M^{-1})$  above 20 K obeys the Curie–Weiss law, with a Curie<br>constant of  $C = 4.47$  amu K/mol and a Weiss constant of  $\theta$ constant of  $C = 4.47$  emu · K/mol and a Weiss constant of  $\theta$  $= -9.7$  K. The *C* value is consistent with the value of 4.37 emu · K/mol expected for a magnetically isolated high-spin Mn<sup>II</sup> ion considering  $g = 2.00$ , while the negative  $\theta$  value suggests a dominant antiferromagnetic coupling among the Mn<sup>II</sup> ions. At an applied field of 5000 Oe,  $\gamma_M T$  decreases slowly as the temperature is decreased to 60 K and then starts to decrease sharply to 2 K. A peak at ca. 6.0 K in the  $\chi_{\text{M}}$ -*T* plot is characteristic for phase transition to an antiferromagnetically ordered structure. However, the magnetic behavior of **2** is more complex in low fields. As shown in Figure 9, upon exposure to a magnetic field of 50 Oe, the value of  $\chi_M T$  first decreases to a minimum at ca. 6 K as the T decreases and then increases to reach a maximum at 4.3 K and finally decreases upon further cooling. Interestingly, the magnetic behavior of **2** is highly dependent on the applied field, and the maximum around 4.3 K in the  $\chi_M T - T$  curve goes down as the field increases. These features are typical for spin-canted antiferromagnetism, $^{22}$  in which a predominantly antiferromagnetic phase possesses a spontaneous

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**Figure 9.** Plots of  $\chi_M T$  vs *T* at different fields for 2. Inset: plots of  $\chi_M$  vs *T* under various fields.

magnetization originating from a small deviation between the antiparallel alignment of the spins. The disappearance of the peak at ca. 6 K in the  $\chi_M-T$  plot at a higher applied field (Figure 9, inset) suggests the occurrence of a fieldinduced spin-flop transition. $^{23}$  The ac susceptibility (Figure S4b in the Supporting Information) shows a maximum at 6 K in the in-phase  $(\chi')$  region, and a peak is also detectable in the out-of-phase  $(\chi'')$  region at the same temperature. The nonzero signal in the imaginary component indicates the presence of a weak ferromagnetism as a result of the spin canting in this magnetic system. Furthermore, the absence of a frequency dependence of ac susceptibility precludes the possibility of the spin glass in **2**.

At 2 K, the field dependence of the magnetization increases almost linearly to the highest field measured of 50 kOe (Figure S5b in the Supporting Information), and no obvious sigmoid-shaped curve is observed. The magnetization value is 2.25  $N\beta$  at 50 kOe, far from the saturation value of 5  $N\beta$ for a  $Mn^{II}$  ion  $(S = 5/2)$ , which is in agreement with an antiferromagnet. In addition a hypersection is observed antiferromagnet. In addition, a hysteresis loop is observed at 2 K with a coercive field  $(H_c)$  of 244 Oe and a remnant magnetization  $(M_r)$  of 0.014  $N\beta$  (Figure 10). This result further suggests the existence of spontaneous magnetization in **2**, being consistent with the spin-canted antiferromagnetic behavior. The canting angle is calculated to be ca. 2.72(3)°, based on the equation  $\theta = \tan^{-1}(M_r/M_s)^{3.24}$  The critical field<br>of the spin-flop transition is determined by a clear disconof the spin-flop transition is determined by a clear discontinuity in the d*M*/d*H* curve (Figure S5b, inset, in the Supporting Information). The appearance of a maximum around 12.3 kOe in the  $\chi'$  vs *H* curve measured at 2 K further confirms the presence of the magnetic phase transition for **2** (Figure 11).

Generally, spin canting can arise from two contributions: (1) the presence of an antisymmetric exchange and (2) the



**Figure 10.** Hysteresis loop for **2** at 2.0 K.



**Figure 11.** Field-dependent ac magnetic susceptibilities of **2** measured at 2 K.

existence of single-ion magnetic anisotropy.<sup>25</sup> It is wellknown that the presence of an inversion center between adjacent spin centers can result in the disappearance of the antisymmetric exchange.<sup>26</sup> However, several high-spin manganese(II) complexes with inversion centers exhibiting spincanted behavior have been reported in spite of the lack of antisymmetric exchange and the isotropic  $Mn^{\text{II}}$  ion, which was regarded as likely arising from weak single-ion anisotropy of the  $Mn^{\text{II}}$  site or anisotropy of the whole crystal to be considered.27 In **1**, spin canting is not observed because of the existence of an inversion center in the  $Mn_2(\mu_2$ -OH)<sub>2</sub> core, while in  $2$ , although the adjacent spin  $Mn^{\text{II}}$  ions are related to a crystallographic inversion center, all of the  $Mn<sub>2</sub>(bta)<sub>2</sub>$  dimers are connected by  $2<sub>1</sub>$  helices with opposite chirality. The different orientations of opposite chirality may induce the occurrence of antisymmetric exchange, thus resulting in spin canting. Similar phenomena have been observed in several manganese azide complexes with the

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same space group,  $28$  in which a screw chain with opposite chirality between neighboring Mn<sub>2</sub> subunits also leads to spin canting. In light of Robison's expression  $\sigma^2 = \sum_{i=1}^{12} (\theta_i -$ 90)<sup>2</sup>/11, where  $\theta_i$  are the bond angles of the octahedron,<sup>29</sup> the square difference of the bond angles  $(\sigma^2)$  of the Mn1<sup>II</sup> coordination octahedron in **2** is 119.3, which shows a clear and great distortion of the  $Mn1<sup>H</sup>$  octahedron. Therefore, local single-ion anisotropy of  $Mn1<sup>H</sup>$  originating from the distorted octahedron enhances the antisymmetric interaction,<sup>30</sup> which might also contribute to spin canting.

In summary, the diversities of deprotonated types and coordination modes of the H2bta ligand generate different structural networks and further lead to distinct magnetic properties. The 2-D gridlike layer of **1** behaves as antiferromagnetic ordering with  $T_N = 4.5$  K, while the 3-D PtS- like framework of **2** displays the coexistence of spin canting and spin-flop transition. These studies show that tetrazolate bridging ligands are particularly promising for the design of extended networks exhibiting different and interesting bulk magnetic behaviors. The research on other metal complexes of the H2bta ligand and their magnetic studies are in process.

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**Supporting Information Available:** TGA plots, PXRD patterns, ac susceptibility measurements, field dependence of magnetization, and IR spectra for 1 and 2, view of the  $-ABAB-$  arrangement for **1**, and X-ray crystallographic data in CIF format. This material is available free of charge via the Internet at http://pubs.acs.org.

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